

Quantum states are relative to observers!

... and this is important for
quantum networks

LINCS reading group “Network Theory”

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The Nokia Bell Labs logo is displayed in white text within a dark blue circular area. The text is arranged in three lines: "NOKIA" on the top line, "BELL" on the middle line, and "LABS" on the bottom line. The letters are in a clean, sans-serif font.

Quantum states are relative to observers!

Pseudo-paradoxes in quantum physics

Some pseudo-paradoxes in quantum physics

- Schrödinger's cat: dead and alive at the same time?
- **Einstein-Podolsky-Rosen (EPR) dilemma**: spooky action at distance?
 - Measuring an entangled qubit in one location instantaneously changes its state of the other qubit in another location

Misunderstanding of what is a quantum state leads to wrong statements

- Special relativity is wrong because of the EPR “paradox”?
 - Special relativity (and general relativity) has always been true since 1905 (resp., since 1916)...
- Transmission of “information” faster than light with entanglement?
 - In contradiction with the no-go theorem about “no-faster-than-light signaling” in quantum physics...
 - Quantum physics has always been true since years 1920s...
- The so-called “measurement problem” in quantum physics?
 - Still an open question in quantum physics, but a better understanding of what is a quantum state should help...

What is a quantum state?

Quantum states are relative to observers!

Some mathematical notations

Dirac notation (“bra-ket” notation)

- See https://fr.wikipedia.org/wiki/Notation_bra-ket
- Quantum state ψ = line (ray) in Hilbert space $V = \mathbb{C}^d$ (N qubits: $d = 2^N$)
- Vector representing quantum state ψ : $|\psi\rangle \in V$ (“ket”)
- Corresponding “bra”: $\langle\psi| \in \text{dual}(V)$
 - Rule: $\lambda|\varphi\rangle + \mu|\psi\rangle \leftrightarrow \lambda^*\langle\varphi| + \mu^*\langle\psi|$
- Hermitian product (\sim scalar product in Euclidian space): $\langle\varphi|\psi\rangle = \langle\varphi| \cdot |\psi\rangle = \langle\psi|\varphi\rangle^*$
$$\langle (\lambda|\varphi\rangle + \mu|\psi\rangle) | (\lambda'|\varphi'\rangle + \mu'|\psi'\rangle) \rangle = (\lambda^*\langle\varphi| + \mu^*\langle\psi|) \cdot (\lambda'|\varphi'\rangle + \mu'|\psi'\rangle)$$
$$= \lambda^*\lambda'\langle\varphi|\varphi'\rangle + \lambda^*\mu'\langle\varphi|\psi'\rangle + \mu^*\lambda'\langle\psi|\varphi'\rangle + \mu^*\mu'\langle\psi|\psi'\rangle$$
- Linear operators: $\sum c_k |\psi_k\rangle\langle\varphi_k|$, with $|\psi_k\rangle\langle\varphi_k|\chi\rangle = |\psi_k\rangle\langle\varphi_k|\chi\rangle = \langle\varphi_k|\chi\rangle \cdot |\psi_k\rangle$
- Orthogonal projector on a state ψ : $|\psi\rangle\langle\psi|$ with $\langle\psi|\psi\rangle = 1$
 - $|\psi\rangle\langle\psi||\psi\rangle = \langle\psi|\psi\rangle \cdot |\psi\rangle = |\psi\rangle$
 - $\forall |\varphi\rangle \in V$ such as $\langle\psi|\varphi\rangle = 0$, i.e., $|\varphi\rangle \perp |\psi\rangle$, $|\psi\rangle\langle\psi||\varphi\rangle = \langle\psi|\varphi\rangle \cdot |\psi\rangle = 0 \cdot |\psi\rangle = 0$

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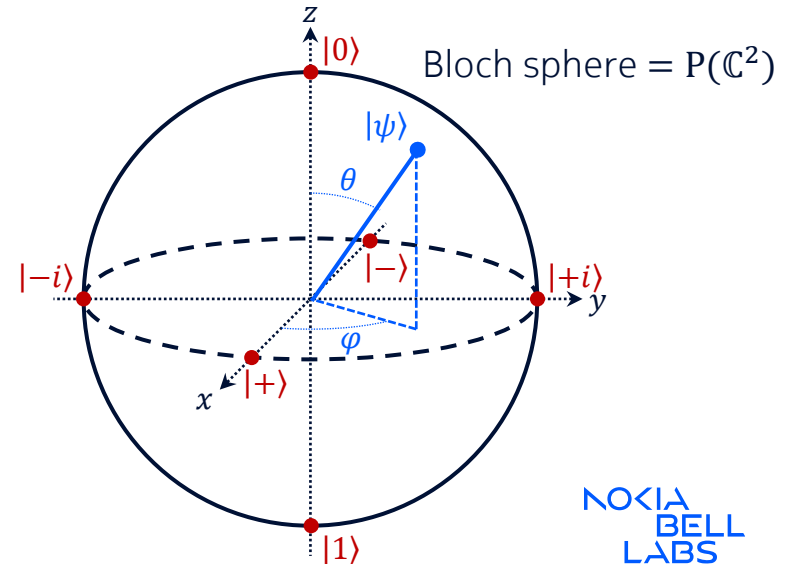
Qubits: bases, Bloch sphere

States in quantum physics (mathematical modeling)

- A quantum state ψ can be represented by a vector in a complex Hilbert space: $|\psi\rangle \in V$ (“ket”), up to normalization factor $\langle\varphi|\varphi\rangle = 1$ and a global phase factor, i.e., $|\varphi\rangle \cong e^{i\alpha}|\varphi\rangle$
- $\psi =$ line (ray) in the Hilbert space: the set of quantum states is the projective space $P(V)$

Qubits = lines in a Hilbert space V with $\dim(V) = 2$

- Quantum state ψ : $|\psi\rangle = (e^{i\alpha}) \cdot \left(\cos\frac{\theta}{2}|0\rangle + e^{i\varphi}\sin\frac{\theta}{2}|1\rangle\right)$
with $\theta \in [0, \pi]$ and $\varphi \in [0, 2\pi[$
- Some bases (orthonormal bases):
 - Computational basis ($|0\rangle, |1\rangle$) on the z axis
 - Basis ($|+\rangle, |-\rangle$) on the x axis with
 $|+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$ and $|-\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$
 - Basis ($|+i\rangle, |-i\rangle$) on the y axis with
 $|+i\rangle = \frac{1}{\sqrt{2}}(|0\rangle + i|1\rangle)$ and $|-i\rangle = \frac{1}{\sqrt{2}}(|0\rangle - i|1\rangle)$



Quantum states are relative to observers!

Qubits: orthogonal state

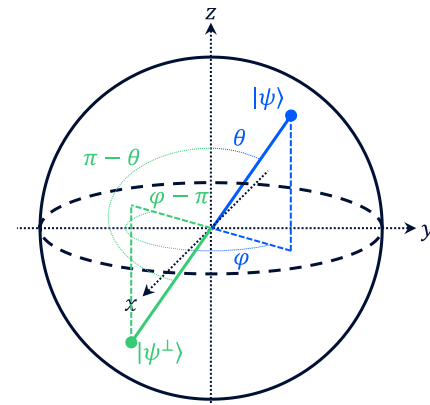
Qubit state ψ represented by the vector $|\psi\rangle = \cos\frac{\theta}{2}|0\rangle + e^{i\varphi}\sin\frac{\theta}{2}|1\rangle$

Its “orthogonal state” ψ^\perp can be represented by the vector:

$$|\psi^\perp\rangle = e^{i\beta} \cdot \left(-e^{-i\varphi}\sin\frac{\theta}{2}|0\rangle + \cos\frac{\theta}{2}|1\rangle \right) = e^{i(\beta-\varphi+\pi)} \cdot \left(\cos\frac{\pi-\theta}{2}|0\rangle + e^{i(\varphi-\pi)}\sin\frac{\pi-\theta}{2}|1\rangle \right)$$

such as $(|\psi\rangle, |\psi^\perp\rangle)$ is another basis of V :

- Computational basis : $\langle 0|0\rangle = \langle 1|1\rangle = 1$ and $\langle 0|1\rangle = \langle 1|0\rangle = 0$
- $\langle \psi|\psi\rangle = \left\langle \cos\frac{\theta}{2}|0\rangle + e^{i\varphi}\sin\frac{\theta}{2}|1\rangle \left| \cos\frac{\theta}{2}|0\rangle + e^{i\varphi}\sin\frac{\theta}{2}|1\rangle \right\rangle = \cos^2\frac{\theta}{2} + e^{-i\varphi}e^{i\varphi}\sin^2\frac{\theta}{2} = 1$
- $\langle \psi^\perp|\psi^\perp\rangle = e^{-i\beta}e^{i\beta} \left\langle -e^{-i\varphi}\sin\frac{\theta}{2}|0\rangle + \cos\frac{\theta}{2}|1\rangle \left| -e^{-i\varphi}\sin\frac{\theta}{2}|0\rangle + \cos\frac{\theta}{2}|1\rangle \right\rangle = e^{+i\varphi}e^{-i\varphi}\sin^2\frac{\theta}{2} + \cos^2\frac{\theta}{2} = 1$
- $\langle \psi^\perp|\psi\rangle = \left\langle -e^{-i\varphi}\sin\frac{\theta}{2}|0\rangle + \cos\frac{\theta}{2}|1\rangle \left| \cos\frac{\theta}{2}|0\rangle + e^{i\varphi}\sin\frac{\theta}{2}|1\rangle \right\rangle = -e^{+i\varphi}\sin\frac{\theta}{2}\cos\frac{\theta}{2} + \cos\frac{\theta}{2}e^{i\varphi}\sin\frac{\theta}{2} = 0$



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Qubits: orthogonal state – Basis transformations

Computational basis ($|\psi\rangle, |\psi^\perp\rangle$) in the basis ($|0\rangle, |1\rangle$):

$$|\psi\rangle = \cos\frac{\theta}{2}|0\rangle + e^{i\varphi}\sin\frac{\theta}{2}|1\rangle \quad (1)$$

$$|\psi^\perp\rangle = e^{i\beta} \cdot \left(-e^{-i\varphi}\sin\frac{\theta}{2}|0\rangle + \cos\frac{\theta}{2}|1\rangle\right) \quad (2)$$

Computational basis ($|0\rangle, |1\rangle$) in the basis ($|\psi\rangle, |\psi^\perp\rangle$):

$$\left(\cos\frac{\theta}{2}\right) \times (1) \Rightarrow \cos\frac{\theta}{2}|\psi\rangle = \cos^2\frac{\theta}{2}|0\rangle + \cos\frac{\theta}{2}e^{i\varphi}\sin\frac{\theta}{2}|1\rangle \quad (3)$$

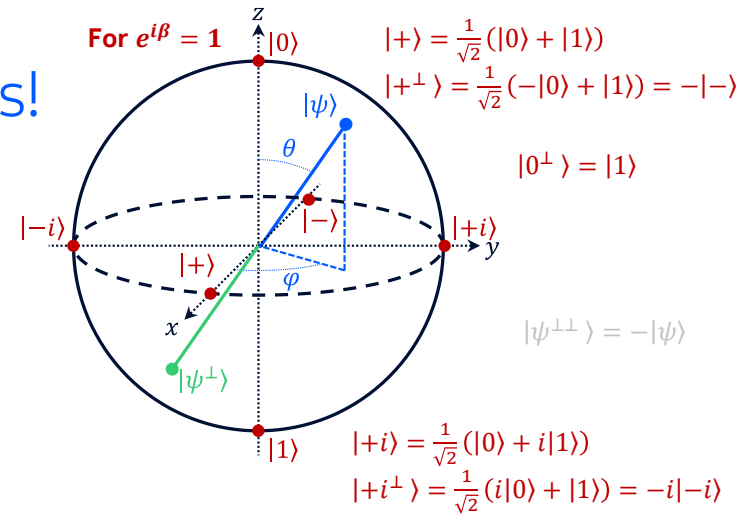
$$\left(-e^{-i\beta}e^{i\varphi}\sin\frac{\theta}{2}\right) \times (2) \Rightarrow -e^{-i\beta}e^{i\varphi}\sin\frac{\theta}{2}|\psi^\perp\rangle = \sin^2\frac{\theta}{2}|0\rangle - e^{i\varphi}\sin\frac{\theta}{2}\cos\frac{\theta}{2}|1\rangle \quad (4)$$

$$(3) + (4) \Rightarrow |0\rangle = \cos\frac{\theta}{2}|\psi\rangle - e^{-i\beta}e^{i\varphi}\sin\frac{\theta}{2}|\psi^\perp\rangle \quad (1')$$

$$\left(e^{-i\varphi}\sin\frac{\theta}{2}\right) \times (1) \Rightarrow e^{-i\varphi}\sin\frac{\theta}{2}|\psi\rangle = e^{-i\varphi}\sin\frac{\theta}{2}\cos\frac{\theta}{2}|0\rangle + \sin^2\frac{\theta}{2}|1\rangle \quad (5)$$

$$\left(e^{-i\beta}\cos\frac{\theta}{2}\right) \times (2) \Rightarrow e^{-i\beta}\cos\frac{\theta}{2}|\psi^\perp\rangle = -\cos\frac{\theta}{2}e^{-i\varphi}\sin\frac{\theta}{2}|0\rangle + \cos^2\frac{\theta}{2}|1\rangle \quad (6)$$

$$(5) + (6) \Rightarrow |1\rangle = e^{-i\varphi}\sin\frac{\theta}{2}|\psi\rangle + e^{-i\beta}\cos\frac{\theta}{2}|\psi^\perp\rangle \quad (2')$$



Quantum states are relative to observers!

Qubits: observables, measurement

Observables on qubits:

- $\forall c \in V, M_{\psi, \lambda_{\psi}, \lambda_{\psi^{\perp}}} = \lambda_{\psi} |\psi\rangle\langle\psi| + \lambda_{\psi^{\perp}} |\psi^{\perp}\rangle\langle\psi^{\perp}|$ with $\lambda_{\psi} \in \mathbb{R}$ and $\lambda_{\psi^{\perp}} \in \mathbb{R} \setminus \{\lambda_{\psi}\}$
- Spin observable: $S_{\psi} = \frac{\hbar}{2} (|\psi\rangle\langle\psi| - |\psi^{\perp}\rangle\langle\psi^{\perp}|)$

Measurement with observable $M_{\psi, \lambda_{\psi}, \lambda_{\psi^{\perp}}}$

- State before measurement: $|\chi\rangle = \cos\frac{\theta}{2} |\psi\rangle + e^{i\varphi} \sin\frac{\theta}{2} |\psi^{\perp}\rangle$
- Measured values and final states:
 - λ_{ψ} with probability $p_{\chi \rightarrow \psi} = \langle\chi|\psi\rangle\langle\psi|\chi\rangle = \cos^2\frac{\theta}{2}$ and final state $|\psi\rangle$
 - With $|\chi \rightarrow \psi\rangle = |\psi\rangle\langle\psi|\chi\rangle = \langle\psi|\chi\rangle|\psi\rangle$, $p_{\chi \rightarrow \psi} = \langle\chi \rightarrow \psi|\chi \rightarrow \psi\rangle$ and final state $\frac{1}{\sqrt{p_{\chi \rightarrow \psi}}} |\chi \rightarrow \psi\rangle$
 - $\lambda_{\psi^{\perp}}$ with probability $p_{\chi \rightarrow \psi^{\perp}} = \langle\chi|\psi^{\perp}\rangle\langle\psi^{\perp}|\chi\rangle = \sin^2\frac{\theta}{2}$ and final state $|\psi^{\perp}\rangle$
 - With $|\chi \rightarrow \psi^{\perp}\rangle = |\psi^{\perp}\rangle\langle\psi^{\perp}|\chi\rangle = \langle\psi^{\perp}|\chi\rangle|\psi^{\perp}\rangle$, $p_{\chi \rightarrow \psi^{\perp}} = \langle\chi \rightarrow \psi^{\perp}|\chi \rightarrow \psi^{\perp}\rangle$ and final state $\frac{1}{\sqrt{p_{\chi \rightarrow \psi^{\perp}}}} |\chi \rightarrow \psi^{\perp}\rangle$

Quantum states are relative to observers!

Pairs of qubits and entanglement: the 4 Bell states

Pair of qubits:

- $\text{Vect}(\mathbb{C}^2 \times \mathbb{C}^2) = \mathbb{C}^2 \otimes \mathbb{C}^2 = \mathbb{C}^4$
- Computational basis: $\{|0\rangle, |1\rangle\} \times \{|0\rangle, |1\rangle\} = \{|0\rangle \otimes |0\rangle, |0\rangle \otimes |1\rangle, |1\rangle \otimes |0\rangle, |1\rangle \otimes |1\rangle\}$
- Some other bases: $\{|\varphi\rangle, |\varphi^\perp\rangle\} \times \{|\psi\rangle, |\psi^\perp\rangle\} = \{|\varphi\rangle \otimes |\psi\rangle, |\varphi\rangle \otimes |\psi^\perp\rangle, |\varphi^\perp\rangle \otimes |\psi\rangle, |\varphi^\perp\rangle \otimes |\psi^\perp\rangle\}$
- Another basis: the 4 Bell states $BSB = \{|\Phi^+\rangle, |\Phi^-\rangle, |\Psi^+\rangle, |\Psi^-\rangle\}$
 - $|\Phi^+\rangle = \frac{1}{\sqrt{2}}(|0\rangle \otimes |0\rangle + |1\rangle \otimes |1\rangle)$
 - $|\Phi^-\rangle = \frac{1}{\sqrt{2}}(|0\rangle \otimes |0\rangle - |1\rangle \otimes |1\rangle)$
 - $|\Psi^+\rangle = \frac{1}{\sqrt{2}}(|0\rangle \otimes |1\rangle + |1\rangle \otimes |0\rangle)$
 - $|\Psi^-\rangle = \frac{1}{\sqrt{2}}(|0\rangle \otimes |1\rangle - |1\rangle \otimes |0\rangle)$
 - One can prove that $\forall \Phi \in BSB, \forall \varphi \in \mathbb{C}^2, \forall \psi \in \mathbb{C}^2, |\Phi\rangle \neq |\varphi\rangle \otimes |\psi\rangle$
 - $\Phi \in BSB$ is a (maximally) entangled state for a pair of qubits

Quantum states are relative to observers!

Entanglement: $|\Psi^-\rangle$ Bell state symmetry

“Central symmetry” of the bell state $|\Psi^-\rangle$

- Bell state $|\Psi^-\rangle$: $|\Psi^-\rangle = \frac{1}{\sqrt{2}} (|0\rangle \otimes |1\rangle - |1\rangle \otimes |0\rangle)$
- Computational basis $(|0\rangle, |1\rangle)$ in the basis $(|\psi\rangle, |\psi^\perp\rangle)$:

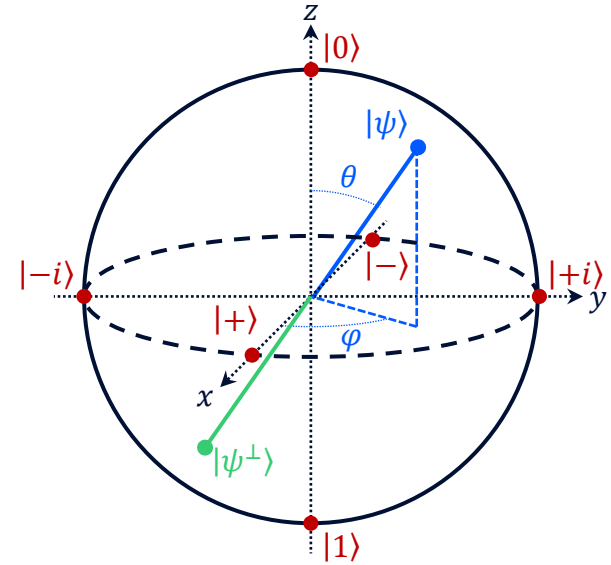
$$|0\rangle = \cos\frac{\theta}{2} |\psi\rangle - e^{-i\beta} e^{i\varphi} \sin\frac{\theta}{2} |\psi^\perp\rangle \quad (1')$$

$$|1\rangle = e^{-i\varphi} \sin\frac{\theta}{2} |\psi\rangle + e^{-i\beta} \cos\frac{\theta}{2} |\psi^\perp\rangle \quad (2')$$

- Computational basis $(|0\rangle, |1\rangle)$ in the basis $(|\psi\rangle, |\psi^\perp\rangle)$:

$$\begin{aligned} |\Psi^-\rangle &= \frac{1}{\sqrt{2}} \left(\cos\frac{\theta}{2} |\psi\rangle - e^{-i\beta} e^{i\varphi} \sin\frac{\theta}{2} |\psi^\perp\rangle \right) \otimes \left(e^{-i\varphi} \sin\frac{\theta}{2} |\psi\rangle + e^{-i\beta} \cos\frac{\theta}{2} |\psi^\perp\rangle \right) \\ &\quad - \frac{1}{\sqrt{2}} \left(e^{-i\varphi} \sin\frac{\theta}{2} |\psi\rangle + e^{-i\beta} \cos\frac{\theta}{2} |\psi^\perp\rangle \right) \otimes \left(\cos\frac{\theta}{2} |\psi\rangle - e^{-i\beta} e^{i\varphi} \sin\frac{\theta}{2} |\psi^\perp\rangle \right) \\ &= 0 \cdot |\psi\rangle \otimes |\psi\rangle + \frac{e^{-i\beta} (\cos^2\frac{\theta}{2} + \sin^2\frac{\theta}{2})}{\sqrt{2}} |\psi\rangle \otimes |\psi^\perp\rangle - \frac{e^{-i\beta} (\sin^2\frac{\theta}{2} + \cos^2\frac{\theta}{2})}{\sqrt{2}} |\psi^\perp\rangle \otimes |\psi\rangle + 0 \cdot |\psi^\perp\rangle \otimes |\psi^\perp\rangle \end{aligned}$$

- Choice $e^{-i\beta} = 1 \Rightarrow \forall \psi \in P(\mathbb{C}^2), |\Psi^-\rangle = \frac{1}{\sqrt{2}} (|\psi\rangle \otimes |\psi^\perp\rangle - |\psi^\perp\rangle \otimes |\psi\rangle)$



Quantum states are relative to observers!

Entanglement: $|\Psi^-\rangle$ Bell state measurements

State before measurement: $|\Psi^-\rangle = \frac{1}{\sqrt{2}} (|\varphi\rangle \otimes |\varphi^\perp\rangle - |\varphi^\perp\rangle \otimes |\varphi\rangle) = \frac{1}{\sqrt{2}} (|\psi\rangle \otimes |\psi^\perp\rangle - |\psi^\perp\rangle \otimes |\psi\rangle)$

Measurement of the 1st qubit in state $\varphi \in \mathcal{P}(\mathbb{C}^2)$:

- $|\Psi^- \rightarrow \varphi \otimes ?\rangle = (|\varphi\rangle\langle\varphi| \otimes Id_2)|\Psi^-\rangle = \frac{1}{\sqrt{2}} (|\varphi\rangle\langle\varphi||\varphi\rangle \otimes |\varphi^\perp\rangle - |\varphi\rangle\langle\varphi||\varphi^\perp\rangle \otimes |\varphi\rangle) = \frac{1}{\sqrt{2}} |\varphi\rangle \otimes |\varphi^\perp\rangle$

- The measured state for the 1st qubit is $|\varphi\rangle$ with probability $\left(\frac{1}{\sqrt{2}}\right)^2 = \frac{1}{2}$
- The state of the pair after measurement is $|\varphi\rangle \otimes |\varphi^\perp\rangle$ (entanglement has disappeared...)
- The inferred state for the 2nd qubit is $|\varphi^\perp\rangle$

Measurement of the 2nd qubit in state $\psi^\perp \in \mathcal{P}(\mathbb{C}^2)$:

- $|\Psi^- \rightarrow ? \otimes \psi^\perp\rangle = (Id_2 \otimes |\psi^\perp\rangle\langle\psi^\perp|)|\Psi^-\rangle = \frac{1}{\sqrt{2}} (|\psi\rangle \otimes |\psi^\perp\rangle\langle\psi^\perp||\psi^\perp\rangle - |\psi^\perp\rangle \otimes |\psi^\perp\rangle\langle\psi^\perp||\psi\rangle) = \frac{1}{\sqrt{2}} |\psi\rangle \otimes |\psi^\perp\rangle$

- The measured state for the 2nd qubit is $|\psi^\perp\rangle$ with probability $\left(\frac{1}{\sqrt{2}}\right)^2 = \frac{1}{2}$
- The state of the pair after measurement is $|\psi\rangle \otimes |\psi^\perp\rangle$ (entanglement has disappeared...)
- The inferred state for the 1st qubit is $|\psi\rangle$

Quantum states are relative to observers!

EPR experiment: description and paradox

EPR experiment

- Pair of qubits in state $|\Psi^-\rangle$
- Alice (observer A) and Bob (observer B) are in different locations
 - Distance $d_{A\leftrightarrow B}$ such as $\Delta t_{A\leftrightarrow B} = \frac{d_{A\leftrightarrow B}}{c}$ is not negligible
- Alice measure the 1st qubit in state $\varphi \in \mathcal{P}(\mathbb{C}^2)$ at time t_A in her inertial frame of reference R
 - Time t'_A in Bob's inertial reference frame R'
 - She received (at the speed of light) Bob's measurement outcome at time $t_{B\rightarrow A} > t_A$
 - Thus, the projected state after her measurement is $|\Psi^- \rightarrow \varphi \otimes ?\rangle_{Alice} = \frac{1}{\sqrt{2}}|\varphi\rangle \otimes |\varphi^\perp\rangle$
- Bob measure the 2nd qubit in state $\psi^\perp \in \mathcal{P}(\mathbb{C}^2) \setminus \{\varphi, \varphi^\perp\}$ at time t'_B in its inertial frame of reference R'
 - Time t_B in Alice's inertial reference frame R
 - He received (at the speed of light) Alice's measurement outcome at time $t'_{A\rightarrow B} > t'_B$
 - Thus, the projected state after his measurement is $|\Psi^- \rightarrow ? \otimes \psi^\perp\rangle_{Bob} = \frac{1}{\sqrt{2}}|\psi\rangle \otimes |\psi^\perp\rangle$
- Paradox = The states after their respective measurements are different:
 $|\varphi\rangle \otimes |\varphi^\perp\rangle \neq |\psi\rangle \otimes |\psi^\perp\rangle \Rightarrow$ Who is right, Alice or Bob?

Quantum states are relative to observers!

EPR experiment: observer-dependent time \Rightarrow observer-dependent state

Usual interpretation: Alice (or Bob) measures first and instantaneous change of remote qubit

- Special Relativity \Rightarrow time depends on the observers' inertial frames of reference R and R'
- We can choose R (resp. R') such as $t_B > t_A$, $t_B < t_A$ or even $t_B = t_A$ (resp. $t'_B > t'_A$, $t'_B < t'_A$ or $t'_B = t'_A$)
- This implies the observer-dependence of the state!

Quantum physics formalism does not care about the order of separated measurements:

$$(Id_2 \otimes |\psi^\perp\rangle\langle\psi^\perp|) \times (|\varphi\rangle\langle\varphi| \otimes Id_2) = (|\varphi\rangle\langle\varphi| \otimes Id_2) \times (Id_2 \otimes |\psi^\perp\rangle\langle\psi^\perp|) = (|\varphi\rangle\langle\varphi| \otimes |\psi^\perp\rangle\langle\psi^\perp|)$$

- The projected state of the pair of qubits for Alice after $t_{B \rightarrow A}$ is:

$$|\Psi^- \rightarrow \varphi \otimes \psi^\perp\rangle_{Alice} = (Id_2 \otimes |\psi^\perp\rangle\langle\psi^\perp|) \left(\frac{1}{\sqrt{2}} |\varphi\rangle \otimes |\varphi^\perp\rangle \right) = \frac{\langle\psi^\perp|\varphi^\perp\rangle}{\sqrt{2}} |\varphi\rangle \otimes |\psi^\perp\rangle$$

- The projected state of the pair of qubits for Bob after $t'_{A \rightarrow B}$ is:

$$|\Psi^- \rightarrow \varphi \otimes \psi^\perp\rangle_{Bob} = (|\varphi\rangle\langle\varphi| \otimes Id_2) \left(\frac{1}{\sqrt{2}} |\psi\rangle \otimes |\psi^\perp\rangle \right) = \frac{\langle\varphi|\psi\rangle}{\sqrt{2}} |\varphi\rangle \otimes |\psi^\perp\rangle$$

$$\Rightarrow |\Psi^- \rightarrow \varphi \otimes \psi^\perp\rangle_{Alice} = |\Psi^- \rightarrow \varphi \otimes \psi^\perp\rangle_{Bob} = \frac{\langle\varphi|\psi\rangle}{\sqrt{2}} |\varphi\rangle \otimes |\psi^\perp\rangle$$

because one can check that $\langle\psi^\perp|\varphi^\perp\rangle = \langle\varphi|\psi\rangle$

Quantum states are relative to observers!

EPR experiment: summary

EPR experiment + Special Relativity \Rightarrow Quantum states may be relative to the observer!

- Before the measurement: same state

$$|initial\ state\rangle_{Alice} = |initial\ state\rangle_{Bob} = |\Psi^-\rangle$$

- After local measurement and before reception of other measurement outcome: different states!

$$|\Psi^- \rightarrow \varphi \otimes ?\rangle_{Alice} = \frac{1}{\sqrt{2}}|\varphi\rangle \otimes |\varphi^\perp\rangle \neq |\Psi^- \rightarrow ? \otimes \psi^\perp\rangle_{Bob} = \frac{1}{\sqrt{2}}|\psi\rangle \otimes |\psi^\perp\rangle$$

$$\text{with } \psi \in P(\mathbb{C}^2) \setminus \{\varphi, \varphi^\perp\}$$

- After reception of other measurement outcome: same state again

$$|\Psi^- \rightarrow \varphi \otimes \psi^\perp\rangle_{Alice} = |\Psi^- \rightarrow \varphi \otimes \psi^\perp\rangle_{Bob} = \frac{\langle \varphi | \psi \rangle}{\sqrt{2}} |\varphi\rangle \otimes |\psi^\perp\rangle$$

- See Quirk circuit <https://www.ludovic-noirie.fr/QC/div/MeasurementPsi-.htm>

Quantum states are relative to observers!

EPR experiment: relational interpretation of quantum physics

Carlo Rovelli's relational interpretation of quantum physics

- Carlo Rovelli, *Relational Quantum Mechanics*, 1996, <https://arxiv.org/abs/quant-ph/9609002>
- See also https://en.wikipedia.org/wiki/Interpretations_of_quantum_mechanics#Relational_quantum_mechanics

“Quantum mechanics is a theory about the physical description of physical systems relative to other systems, and this is a complete description of the world”

- Quantum state = mathematical modeling of the state of knowledge of the observer on the observed system
- Quantum states are relative to observers (like time in special and general relativity)
- All systems are quantum systems (main difference with Copenhagen interpretation)
- Observation = Physical interaction = Entanglement (observer out) or Measurement (observer in)

View of Asher Peres (one of the inventors of quantum teleportation)

- Asher Peres, *Quantum information and general relativity*, 2004, <https://arxiv.org/abs/quant-ph/0405127>
- See also https://en.wikipedia.org/wiki/Asher_Peres#Views_on_the_EPR_paradox
- Quantum state = information
- When Alice measures her qubit, absolutely nothing happens at Bob's location
- Bob needs information from Alice to change his states about his qubit

Quantum states are relative to observers!

Schrödinger's cat, interactions and measurement in quantum physics

The “observed observer” (*Relational Quantum Mechanics*, Carlo Rovelli)

- Schrödinger's cat thought experiment: cat in a box in a superposition state “dead or/and alive”
 - System S observed by observer A inside the box
 - Same system S observed by observer B outside the box
- The state is relative to the observer in this configuration too:
 - Observer A: The cat (system S) is dead $|0\rangle_S$ or alive $|1\rangle_S$, not both...
 - Observer B: The cat and the observer A are entangled in superposition state $\frac{1}{\sqrt{2}}(|0\rangle_A|0\rangle_S + |1\rangle_A|1\rangle_S)$
- Is there really a “measurement problem” in quantum physics?
 - Observation = Acquisition of information about a system by another system (by interaction)
 - **Measurement** = Interaction between two systems (S and A), the observer being one of them (A)
 - **Entanglement** = Interaction between two systems (S and A) observed by another system (B)
- Qubit = smallest measurement apparatus to measure another qubit: see Quirk circuit
<https://www.ludovic-noirie.fr/QC/div/MeasurementQubit.htm>

Quantum states are relative to observers!

Conclusion

Quantum states may be relative to observers

- Case 1 (EPR-like): two distant observers on the same composite quantum system with two entangled components, each observer measuring one component and each component being “significantly” distant from the other
- Case 2 (the “observed observer”): a 1st observer observing a quantum system, a second observer observing the 1st observer and the observed quantum system (measurement vs. entanglement)

When is it relevant for case 1?

- Not relevant for quantum computers: single observer, qubits in the “same location” (30 cm = 1 ns)
- But clearly relevant in quantum distributed systems such as quantum networks / quantum internet
 - Each node observes a qubit entangled in Bell state with another qubit of another node
 - The nodes are “significantly” distant (photon transmission in fiber: 50 km = 0.25 ms)
- This effect is important for QKD application in quantum internet: statistical detection of eavesdropper
 - See <https://www.ludovic-noirie.fr/QC/div/QKD.htm>

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Bloch spheres

